

*Tube derivative for non-cylindrical problems
Application to Shape-Newton Method*

Raja Dziri — Jean Paul Zolésio

N° 4676

Décembre 2002

THÈME 4



*apport
de recherche*

Tube derivative for non-cylindrical problems Application to Shape-Newton Method

Raja Dziri , Jean Paul Zolésio

Thème 4 — Simulation et optimisation
de systèmes complexes
Projet Opale

Rapport de recherche n° 4676 — Décembre 2002 — 26 pages

Résumé : The purpose of this preport is to give the expression of the shape derivative of functionals related to moving domains evolution problems. The main step is to define pertubations of a given non-cylindrical domain [also called tube] using transverse vector fields. Under a smoothness assumption, a tube can be generated by a suitable vector field starting from its initial domain. This relationship tube-vector field allows us to define the field derivative and to recognize the term which constitutes the shape derivative for a class of non-cylindrical functionals. We conclude the paper, by applying the obtained result to compute the second order optimality condition for a shape minimization problem.

Mots-clés : Non-cylindrical shape functionals, Eulerian derivatives, transverse vector fields, shape gradient and hessian

Dérivée de forme pour des problèmes non-cylindriques

Abstract: Le but de ce rapport est de donner l'expression de la dérivée de forme de fonctionnelles associées à des problèmes d'évolution non-cylindrique. Le point clé de ce travail réside dans la définition des perturbations d'un domaine non-cylindrique [appelé aussi tube] en utilisant des champs de vecteurs transverses. Sous des hypothèses de régularités, un tube peut être généré par un champ de vecteur convenable connaissant son domaine initial. Cette relation tube-champ de vecteur nous permet de définir une dérivée par rapport aux champs de vitesse et d'en déduire l'expression de la dérivée par rapport à la forme. On analyse la dérivation de fonctionnelle de tube associées à des problèmes dynamiques non cylindriques. On considère le cas particulier de l'évolution d'une fonctionnelle de domaine et l'application aux méthode d'ordre 2 de type Newton; conclut ce rapport, en donnant la condition nécessaire s d'optimalité du second ordre(faisant intervenir le hessien de forme) pour un tel problème de minimisation.

Key-words: Fonctionnelles de forme non-cylindrique, dérivation par rapport aux champs de vitesse, champ transverse, dérivation par rapport à la forme

Table des matières

1	Introduction	4
2	Basic tools	6
2.1	Transformations-Velocity method	6
2.2	Tube perturbation	7
2.3	Explicit expression of $\mathbf{Z}(t).n_{\Omega_t(V)}$	12
2.4	Adjoint problem associated to \mathbf{Z}	13
2.5	A right-hand term supported by $\Sigma(V)$	14
3	Derivability with respect to the field	17
4	Newton-Shape Method	21
4.1	expansion	22
4.2	Second derivative with density gradient formulation	23
4.3	Example	25

1 Introduction

Dynamical shape control problems for systems modeled with dynamical non-cylindrical Partial Differential Equations (PDE) are encountered in fluid-structure issues, free boundary problems, etc...

The cost functional involved in such problems is expressed in terms of integrals over the non-cylindrical evolution domain and/or its lateral boundary. Following [] this non-cylindrical evolution domain will be called a tube and is of the following form :

$$Q = \bigcup_{0 < t < \tau} (\{t\} \times \Omega_t); \quad \text{At } t = 0, \Omega_0 = \Omega.$$

The initial geometry Ω is called the base of the tube. Non smooth tubes are described in [11] with help of the time convection of the base by non lipschitzian vector field V . If the lateral boundary Σ of the tube Q is smooth enough to guarantee the existence of the outward normal field $\vec{\nu}$ to Σ , there exists smooth non-autonomous vector fields V such that

$$T_t(V)\Omega = \Omega_t \subset R^N \quad \forall t \in [0, \tau] \quad (1)$$

where $T(V)$ is the flow associated to V . More precisely, the *intrinsic* outward normal field $\vec{\nu} \in R^{N+1}$ defined on the lateral boundary Σ of Q , can be written as (cf.[9])

$$\vec{\nu}(t) = \frac{1}{\sqrt{1 + v_\nu^2}} (-v_\nu(t), \vec{n}_{\Omega_t})$$

where $\vec{n}_{\Omega_t} \in R^N$ is the “horizontal” outward normal field to Ω_t . Then any sufficiently smooth non-autonomous vector field V such that

$$V(t) \cdot \vec{n}_{\Omega_t} = v_\nu(t) \quad \text{on } \Sigma \quad (2)$$

builds the tube Q (or equivalently : satisfies (1)).

Conversely, to any non-autonomous vector field V , one can associate, in the time interval $[0, \tau]$, a tube $Q(V)$ (also denoted Q_V) by setting $\Omega_t = T_t(V)(\Omega)$.

Obviously (2) is then satisfied on the lateral boundary $\Sigma(V)$ of $Q(V)$.

The functionals considered may depend not only on the tube containing the evolution takes place but also on a field that builds this tube which verifies some extra physical condition (cf. [6]). In that specific situation, the functional is no more a tube functional as it depends on the tube but also on a specific vector field V . For that reason we consider now as independant variable in the problem, not the tube itself but the vector field V whose flow mapping builds the tube starting from a given base Ω . We designate by Q_V the tube built by an admissible vector field V and we consider functionals $j(V)$ in the form

$$j(V) = J(V, Q_V).$$

When J depends exclusively on Q_V we get :

$$j(V + W) = j(V), \quad \forall W \text{ s.t. } \langle W, n_{\Omega_t(V)} \rangle = 0 \quad \text{on } \Sigma_V,$$

this means that \mathbf{j} depends only on the shape of the tube (it is, thus, called a tube functional). The dependence of J on a field V comes, generally, from boundary conditions associated to the state equation.

For a bounded tube Q , there exists a bounded open set D (called a hold-all) such that Q and its perturbations remain in $(0, \tau) \times D$. That cylindrical tube will be invariant under all the transformations we shall consider so that the vector field V as well as the perturbations fields W will satisfy the condition :

$$\langle W(t), \vec{n}_D \rangle = 0 \quad \text{on } (0, \tau) \times \partial D \quad (3)$$

where \vec{n}_D is the outward normal field to D , cf. for example [8] or [9].

The aim is to characterize the derivative with respect to the field of the functional $\mathbf{j}'(\mathbf{V}, \mathbf{W})$. In a first step we obtain a general expression in terms of the *transverse* field Z . We study the problem whose Z is solution and we explicit the gradient with use of the transposed equation solution Λ .

We give sufficient conditions under which \mathbf{j} is Gâteaux differentiable at a field V . For an optimization problem, we prove that if the state depends on the field only by the boundary conditions, the field gradient $G(V)$ is supported by Σ_V . But if it depends only on the shape of the tube, $G(V)$ has the following form

$$G(V)(t) = \gamma_{\Gamma_t(V)}^*(g(V)(t)n_{\Omega_t(V)})$$

where $\Gamma_t(V) = \partial(\Omega_t(V))$ and $\gamma_{\Gamma_t(V)}^*$ is the adjoint of the trace operator on $\Gamma_t(V)$, $g(V)(t) \in [\mathcal{D}^{k-1}(\Gamma_t(V))]'$. For these results one has to assume Ω_t to be \mathcal{C}^k , $\forall t \in [0, \tau]$.

2 Basic tools

2.1 Transformations-Velocity method

First, let us recall some basic results dealing with velocity flows. Let $\tau > 0$, D a convex, smooth and bounded open set in R^N . Consider a vector field

$$V : [0, \tau] \times \overline{D} \longrightarrow R^N, (t, x) \longmapsto V(t)(x) \stackrel{def}{=} V(t, x) \text{ such that}$$

$$V(t).n_D = 0 \text{ on } \partial D, \forall t \in [0, \tau]. \quad (4)$$

Moreover, we assume

$$(V) \left\{ \begin{array}{l} \forall x \in \overline{D}, V(., x) \in C([0, \tau]; R^N) \\ \exists c > 0, \forall x, y \in D, \|V(., x) - V(., y)\|_{C([0, \tau]; R^N)} \leq c|x - y| \end{array} \right.$$

where $V(., x)$ is the function $t \longmapsto V(t, x)$. For any $X \in \overline{D}$, associate the solution $x_V(., X)$ of the ordinary differential equation

$$\left\{ \begin{array}{l} \frac{dx}{dt}(t) = V(t, x(t)), t \in (0, \tau) \\ x(0) = X \end{array} \right.$$

For any $t \in [0, \tau]$, we have a transformation

$$T_t(V) : \overline{D} \longrightarrow \overline{D}; X \longmapsto T_t(V)(X) \stackrel{def}{=} x_V(t, X).$$

The mapping $(t, X) \longmapsto T_t(V)(X)$ is denoted $T(V)$ or T if no confusion is possible.

Under assumption (V), the map T has the following properties :

$$\left\{ \begin{array}{l} (T_1) \quad \forall X \in \overline{D}, \quad T(., X) \in C^1([0, \tau]; R^N) \text{ and } \exists c > 0, \\ \quad \forall X, Y \in \overline{D}, \quad \|T(., X) - T(., Y)\|_{C^1([0, \tau]; R^N)} \leq c|X - Y| \\ (T_2) \quad \forall t \in [0, \tau], \quad X \longmapsto T_t(X) : \overline{D} \longrightarrow \overline{D} \text{ is bijective} \\ (T_3) \quad \forall x \in \overline{D}, \quad T^{-1}(., x) \in C([0, \tau]; R^N) \text{ and } \exists c > 0, \\ \quad \forall x, y \in \overline{D}, \quad \|T^{-1}(., x) - T^{-1}(., y)\|_{C([0, \tau]; R^N)} \leq c|x - y| \end{array} \right.$$

Introduce the Banach space, $k \in N$

$$\mathcal{V}_o^k(D) = \{v \in C^k(\overline{D}, R^N) \mid v.n_D = 0 \text{ on } \partial D\}$$

and the following notations :

$$c(F) \stackrel{def}{=} \sup_{y \neq x} \frac{|F(y) - F(x)|}{|y - x|}; \quad c_k(F) \stackrel{def}{=} \sum_{|\alpha|=k} c(\partial^\alpha F) \text{ for } k \geq 1.$$

The regularity result stated below for the flow associated to a given vector field V is developed in [3].

Proposition 1

For all $V \in C([0, \tau], \mathcal{V}_o^k(D))$, $k \in \mathbb{N}$, such that

$$c_k(V(t)) \leq c \quad (\text{for a constant } c > 0 \text{ independent of } t), \quad (5)$$

it is associated a unique map

$$T(V) \in C^1([0, \tau], C^k(\overline{D}, R^N)) \cap C([0, \tau], W^{k+1, \infty}(D, R^N)).$$

Moreover the mapping $t \longrightarrow V(t) \circ T_t$ is in $L^\infty(0, \tau; W^{k+1, \infty}(D, R^N))$.

The transformation $T_t(V)$ is one-to-one and maps \overline{D} into \overline{D} . For each $t \in [0, \tau]$, one can consider the transformation T_t^{-1} and notice that it is the flow, at $s = t$, of the vector field V_t defined by $V_t(s) = -V(t - s)$ (cf. [8], [3]). Using this argument we can prove that $T^{-1} \in C([0, \tau], C^k(\overline{D}, R^N))$ and with the implicit function theorem one can prove the following regularity result.

Proposition 2

The mapping $[0, \tau] \longrightarrow C^k(\overline{D}, R^N)$, $t \longmapsto T_t^{-1}$ is continuously differentiable.

Proof. Let $t \in (0, \tau)$. The map $T_t \in C^k(\overline{D}, R^N)$. Consider the mapping

$$\Phi : [0, \tau] \times C^{k-1}(\overline{D}, R^N) \longrightarrow C^{k-1}(\overline{D}, R^N)$$

defined by

$$\Phi(t, S) = T(V)(t, S) - Id$$

It is clear that $\Phi(t, T_t(V)^{-1}) = 0$. Moreover Φ is continuously differentiable. The partial derivatives are :

$$\partial_t \Phi(t, S) = V(t, T_t(V) \circ S); \quad \partial_S \Phi(t, S) \xi = (DT_t) \circ S \cdot \xi$$

It is clear that $\partial_S \Phi(t, S)$ is in $\text{aut}(C^{k-1}(\overline{D}, R^N))$. We deduce from the implicit function theorem that $T^{-1} \in C^1([0, \tau]; C^{k-1}(\overline{D}, R^N))$. Then the mapping $t \longrightarrow DT_t^{-1} = (DT_t)^{-1} \circ T_t^{-1}$ is also in $C^1([0, \tau]; C^{k-1}(\overline{D}, R^N))$. Therefore the announced regularity is proved. \square

2.2 Tube perturbation

From now on we assume $k \geq 1$. A perturbation of a tube Q_V in a direction W can be obtained by considering transverse transformations (which are transformations acting on domains built at the same time).

Let $\Omega \subset\subset D$ of class C^k be given, for any sufficiently small positive parameter s , we consider the moving domain $Q_{(V+sW)}$ as the perturbation of the tube Q_V in the direction of the field W . It is composed of the sets

$$\Omega_t(V + sW) = T_t(V + sW)(\Omega), \forall t \in [0, \tau].$$

A transverse transformation is a function which maps $\Omega_t(V)$ onto $\Omega_t(V + sW)$ for any $t \in [0, \tau]$ (and \overline{D} onto \overline{D}). A quite natural one is

$$\mathcal{T}_s^t = T_t(V + sW) \circ T_t(V)^{-1}.$$

If the regularity assumptions $(T_1) - (T_3)$ are satisfied by the mapping $(s, x) \mapsto \mathcal{T}_s^t(x)$ for any $t \in [0, \tau]$, this transformation can be considered as the flow of the vector field [see for instance [1]]

$$\mathcal{Z}^t(s, \cdot) = \left(\frac{\partial}{\partial s} \mathcal{T}_s^t \right) \circ \mathcal{T}_s^t(\cdot)^{-1} = [\partial_s T_t(V + sW)] \circ T_t(V + sW)^{-1}. \quad (6)$$

Lemme 1

Let I_0 be a neighborhood of zero; V and W being in $\mathcal{C}([0, \tau]; \mathcal{V}_o^k(D))$ satisfying condition (5). The mapping

$$\begin{aligned} I_0 &\longrightarrow \mathcal{C}([0, \tau]; C^{k-1}(\overline{D}, R^N)) \\ s &\longrightarrow T(V + sW) \end{aligned}$$

is continuously differentiable and $\partial_s(T_t(V + sW))$ satisfies for any $t \in [0, \tau]$,

$$\begin{aligned} \partial_s[T_t(V + sW)] &= \int_0^t D(V + sW)(\mu, T_\mu(V + sW)) \partial_s[T_\mu(V + sW)] d\mu \\ &+ \int_0^t W(\mu, T_\mu(V + sW)) d\mu. \end{aligned} \quad (7)$$

Proof.

$$\begin{aligned} &T_t(V + sW) - T_t(V + s_o W) \\ &= \int_0^t (V + sW)(\mu, T_\mu(V + sW)) - (V + s_o W)(\mu, T_\mu(V + s_o W)) d\mu. \\ &\|T_t(V + sW) - T_t(V + s_o W)\|_{C^{k-1}(\overline{D})} \leq |s - s_o| \int_0^t \|W(\mu)\|_{C^{k-1}(\overline{D})} d\mu \\ &+ \max_{t \in [0, \tau]} \|DV(t) + s_o DW(t)\|_{C^{k-1}(\overline{D})} \int_0^t \|T_\mu(V + sW) - T_\mu(V + s_o W)\|_{C^{k-1}(\overline{D})} d\mu \end{aligned}$$

Applying the Gronwall inequality, it comes

$$\begin{aligned} & \|T_t(V + sW) - T_t(V + s_o W)\|_{C^{k-1}(\overline{D})} \leq \tau |s - s_o| \|W\|_{C([0, \tau]; C^{k-1}(\overline{D}))} \\ & + \tau |s - s_o| \|W\|_{C([0, \tau]; C^{k-1}(\overline{D}))} \int_0^t \exp(t - \mu) d\mu \end{aligned}$$

Thus for any $t \in [0, \tau]$:

$$\|T_t(V + sW) - T_t(V + s_o W)\|_{C^{k-1}(\overline{D})} \leq \tau |s - s_o| \|W\|_{C([0, \tau]; C^{k-1}(\overline{D}))} e^t.$$

This proves that the considered map is in $W^{1, \infty}(I_0; C([0, \tau]; C^{k-1}(\overline{D}; R^N)))$ and also gives the following uniform boundedness (with respect to s) :

$$\begin{aligned} & \frac{1}{|s - s_o|} \| (V + sW)(\mu, T_\mu(V + sW)) - (V + s_o W)(\mu, T_\mu(V + s_o W)) \|_{C^{k-1}(\overline{D})} \\ & \leq \|W(\mu)\|_{C^{k-1}(\overline{D})} + \tau \exp \tau \max_{t \in [0, \tau]} \|DV(t) + s_o DW(t)\|_{C^{k-1}(\overline{D})} \|W(\mu)\|_{C^{k-1}(\overline{D})} \end{aligned}$$

Then according to the Lebesgue theorem, the derivative exists everywhere in I_0 and satisfies (7). It has the following expression :

$$\begin{aligned} \partial_s [T_t(V + sW)] &= \int_0^t \exp\left\{ \int_\xi^t D(V + sW)(\mu, T_\mu(V + sW)) d\mu \right\} W(\xi, T_\xi(V + sW)) d\xi. \\ &= \int_0^t DT_t(V + sW) \cdot [DT_\xi(V + sW)]^{-1} W(\xi, T_\xi(V + sW)) d\xi. \end{aligned}$$

It is clear that this expression is continuous in I_0 .

Let $\mathcal{S}^t(s) = \partial_s [T_t(V + sW)]$. Then $\mathcal{Z}^t(s, x) = \mathcal{S}^t(s) \circ T_t(V + sW)^{-1}$. In the next section, it will be shown that the field derivatives of non-cylindrical functionals are expressed in terms of the transverse vector field $\mathbf{Z}(t, x) = \mathcal{Z}^t(0, x)$. Therefore, one has to know more about this vector field. In particular, it will be shown that \mathbf{Z} can be characterized as the unique solution of

$$\partial_t \mathbf{Z} + [\mathbf{Z}, V] = W \text{ in } (0, \tau) \times D \quad (8)$$

$$\mathbf{Z}(0, \cdot) = 0 \text{ in } D \quad (9)$$

where $[\cdot, \cdot]$ denotes the Lie Brackets. For that purpose let us consider the vector field S , $S(t, \cdot) \stackrel{def}{=} \mathcal{S}^t(0, \cdot) = \mathbf{Z}(t) \circ T_t(V)$.

Lemme 2

The function S is the unique vector field, in $\mathcal{C}^1([0, \tau]; C^{k-1}(\overline{D}, R^N))$, satisfying

$$S(t) = \int_0^t W(\mu, T_\mu(V)) d\mu + \int_0^t DV(\mu, T_\mu(V)) S(\mu) d\mu. \quad (10)$$

Proof. Let \mathcal{F} be the mapping defined by

$$\begin{aligned}\mathcal{F} : [0, \tau] \times C^{k-1}(\overline{D}, R^N) &\rightarrow C^{k-1}(\overline{D}, R^N) \\ (t, \varphi) &\rightarrow DV(t, T_t(V))\varphi + W(t, T_t(V)).\end{aligned}$$

For any $t \in [0, \tau]$ and any $\varphi \in C^{k-1}(\overline{D})$, $\mathcal{F}(t)$ is affine and $\mathcal{F}(\cdot, \varphi)$ is continuous. So the existence and uniqueness of (10) are given by the Cauchy-Lipschitz theorem. Moreover the solution has the following expression

$$\begin{aligned}S(t) &= \int_0^t \exp\left\{\int_s^t DV(\mu, T_\mu(V)) d\mu\right\} W(s, T_s(V)) ds, \quad \forall t \in [0, \tau]. \\ &= \int_0^t DT_t(V) \cdot [DT_s(V)]^{-1} W(s, T_s(V)) ds\end{aligned}$$

□

Lemme 3

If $\psi \in C^1([0, \tau], C^{k-1}(\overline{D}, R^N))$ is such that

$$\partial_t \psi + D\psi \cdot V \in C([0, \tau], C^{k-1}(\overline{D}, R^N))$$

and satisfies (8)-(9), then $\psi \circ T(V)$ belongs to $C^1([0, \tau], C^{k-1}(\overline{D}, R^N))$ and satisfies (10). Conversely, if $\varphi \in C^1([0, \tau], C^{k-1}(\overline{D}, R^N))$ is solution of (10), then $\varphi \circ T(V)^{-1} \in C^1([0, \tau], C^{k-1}(\overline{D}, R^N))$ such that

$$\partial_t [\varphi \circ T(V)^{-1}] + D[\varphi \circ T(V)^{-1}] \cdot V \text{ is in } C([0, \tau], C^{k-1}(\overline{D}, R^N))$$

and satisfies (8)-(9).

Proof. If ψ belongs to $C^1([0, \tau], C^{k-1}(\overline{D}, R^N))$, then it is the same for the mapping $t \mapsto \psi(t) \circ T_t$ and we have :

$$\partial_t (\psi(t, T_t)) = [DV(t)] \circ T_t \cdot \psi(t, T_t) + W(t) \circ T_t.$$

Thus $\psi \circ T(V)$ satisfies (10). Conversely, since $\varphi \in C^1([0, \tau], C^{k-1}(\overline{D}, R^N))$ is solution of (10) and T^{-1} is in $C^1([0, \tau], C^{k-1}(\overline{D}, R^N))$, it comes that $\varphi \circ T(V)^{-1}$ is in $C^1([0, \tau], C^{k-1}(\overline{D}, R^N))$ and we have :

$$\begin{aligned}\partial_t (\varphi(t, T_t^{-1})) &= [\partial_t \varphi] \circ T_t^{-1} + [D\varphi] \circ T_t^{-1} \cdot \partial_t (T_t^{-1}) \\ &= \{[DV] \circ T_t \cdot \varphi + W \circ T_t\} \circ T_t^{-1} - (D\varphi) \circ T_t^{-1} D(T_t^{-1}) \cdot V(t) \\ &= DV(\varphi \circ T_t^{-1}) + W - D(\varphi \circ T_t^{-1}) \cdot V(t)\end{aligned}$$

which concludes the proof. □

Theorem 1

The field \mathbf{Z} is the unique vector field in

$C^1([0, \tau], C^{k-1}(\overline{D}, R^N))$, such that $\partial_t \mathbf{Z} + D\mathbf{Z}.V \in C([0, \tau], C^{k-1}(\overline{D}, R^N))$, solution of problem (8)-(9).

Proof. Consider the solution S of (10), compose by $T(V)^{-1}$ and make use of the lemma 3. \square

Remark 1

Taking into account the characterization of \mathbf{Z} , we deduce that it has the following expression :

$$\mathbf{Z}(t) = \left\{ \int_0^t DT_t(V) \cdot [DT_s(V)]^{-1} W(s, T_s(V)) ds \right\} \circ T_t(V)^{-1}. \quad (11)$$

By the way a useful property of \mathbf{Z} is that if V and W are of free divergence then \mathbf{Z} too. Indeed

Proposition 3

Let V and W in $C([0, \tau], \mathcal{V}_o^k(D))$, $k \geq 1$, such that (5) holds. Assume

$$\operatorname{div} V = \operatorname{div} W = 0 \text{ in } D.$$

Then the field \mathbf{Z} is of divergence free :

$$\operatorname{div} \mathbf{Z} = 0 \text{ in } D.$$

Proof. Let $f \in \mathcal{D}(D)$. The transformations $T_t(V + sW)$ and $T_t(V)^{-1}$ maps D onto D . Then,

$$\int_D f \circ T_s^t dx = \int_D f dx.$$

Indeed since V and W are of free divergence, we have

$$\int_D f dx = \int_{T_t(V+sW)(T_t(V)^{-1}(D))} f dx = \int_{T_t(V)^{-1}(D)} f \circ T_t(V + sW) dx = \int_D f \circ T_s^t dx.$$

From this we deduce that

$$\frac{d}{ds} \left(\int_D f \circ T_s^t dx \right) = 0$$

The mapping $s \mapsto T_s^t$ is in $C^1(I_0; C^{k-1}(\overline{D}, R^N))$, thus

$$\int_D \nabla f \cdot \mathbf{Z} dx = 0, \forall f \in \mathcal{D}(D).$$

or equivalently $\operatorname{div} \mathbf{Z} = 0$ in $\mathcal{D}'(D)$. \square

2.3 Explicit expression of $\mathbf{Z}(t).n_{\Omega_t(V)}$

In the sequel, we consider a set Ω of class \mathcal{C}^k (k is kept greater than 1). As we will see in the applications, the expression of the Eulerian derivative depends on \mathbf{Z} . Hence, it seems to be necessary to introduce two adjoint states. One associated to the state equation and the other to the field \mathbf{Z} . For a functional defined on a tube Q_V (with the initial domain Ω), this unusual situation might be avoided by considering the function \mathbf{z} defined by

$$\mathbf{z}(t) = (\mathbf{Z}(t).n_t) \circ T_t(V) \quad \text{on } (0, \tau) \times \Gamma \quad (\Gamma = \partial\Omega).$$

Lemme 4

The mapping $t \mapsto T_t(V)$ is in $C^1([0, \tau]; C^k(\overline{D}, R^N))$. Therefore

$$t \mapsto n_t \circ T_t = \frac{*(DT_t)^{-1} n}{||*(DT_t)^{-1} n||} \quad \text{is in } C^1([0, \tau]; C^{k-1}(\Gamma))$$

n and n_t are the outward normal fields respectively to Ω and $\Omega_t(V)$, on Γ and Γ_t . Its derivative is given by :

$$\partial_t(n_t \circ T_t) = \langle DV.n_t, n_t \rangle \circ T_t \, n_t \circ T_t - *DV \circ T_t \, n_t \circ T_t.$$

Proposition 4

The function $\mathbf{z} \in C^1([0, \tau]; C^{k-1}(\Gamma))$ is the unique solution of

$$\partial_t \mathbf{z}(t) - \alpha(t) \circ T_t(V) \mathbf{z}(t) = \beta(t) \circ T_t(V) \quad \text{on } (0, \tau) \times \Gamma \quad (12)$$

$$\mathbf{z}(0) = 0 \quad \text{on } \Gamma \quad (13)$$

where $\alpha(t) = \langle DV.n_t, n_t \rangle$, $\beta(t) = W(t).n_t$.

Proof. The mapping $t \mapsto \mathbf{Z}(t, T_t) = S(t)$ is differentiable. Thus

$$\begin{aligned} \partial_t [\mathbf{Z}(t, T_t).n_t \circ T_t] &= \partial_t (\mathbf{Z} \circ T_t) . n_t \circ T_t + \mathbf{Z}(t, T_t). \partial_t (n_t \circ T_t) \\ &= \langle (W(t) + DV(t).Z(t)) \circ T_t, n_t \circ T_t \rangle + \langle DV.n_t, n_t \rangle \circ T_t (\mathbf{Z}(t).n_t) \circ T_t \\ &\quad - \langle (DV.\mathbf{Z}) \circ T_t, n_t \circ T_t \rangle \\ &= (W(t).n_t) \circ T_t + \langle DV.n_t, n_t \rangle \circ T_t (\mathbf{Z}(t).n_t) \circ T_t. \end{aligned}$$

Eventually the desired result is obtained. □

Remark 2

Expression of \mathbf{z} in terms of the data :

From (12) we deduce that

$$\begin{aligned} \partial_t[\mathbf{z} \exp(\int_0^t \alpha(s) \circ T_s ds)] &= [\alpha(t) \circ T_t \mathbf{z} + \partial_t \mathbf{z}] \exp(\int_0^t \alpha(s) \circ T_s ds) \\ &= \beta(t) \circ T_t \exp(\int_0^t \alpha(s) \circ T_s ds). \end{aligned}$$

Hence \mathbf{z} can be expressed as follows

$$\begin{aligned} \mathbf{z}(t) &= \int_0^t \beta(s) \circ T_s(V) \exp(-\int_s^t \alpha(r) \circ T_r(V) dr) ds. \\ &= \int_0^t [W(s).n_s] \circ T_s(V) \exp(-\int_s^t \langle DV(r)n_r, n_r \rangle \circ T_r(V) dr) ds. \end{aligned} \quad (14)$$

2.4 Adjoint problem associated to \mathbf{Z}

As shown in the proof of Lemma 3 the solution of (8)-(9) is obtained via a change of variable. Then if $H(D)$ is a Banach space of functions defined on D , stable by multiplication by functions in $C^{k-1}(\overline{D})$, the same process generates the solution of the adjoint problem associated to \mathbf{Z} .

Theorem 2

Let $F \in L^2((0, \tau); H(D))$. There exists a unique $\Lambda \in C([0, \tau]; H(D))$ such that $\partial_t \Lambda + D\Lambda.V \in L^2((0, \tau); H(D))$ solution of

$$-\partial_t \Lambda - D\Lambda.V - {}^*DV.\Lambda - (\operatorname{div} V) \Lambda = F \quad (15)$$

$$\Lambda(\tau) = 0. \quad (16)$$

Proof. Consider $\theta \in C^1([0, \tau]; H(D))$ the unique solution of the backward problem

$$\begin{aligned} -\partial_t \theta - [{}^*(DV) \circ T_t + (\operatorname{div} V) \circ T_t] \theta &= F \circ T_t \\ \theta(\tau) &= 0. \end{aligned}$$

Applying $\exp \int_0^t [({}^*DV(s)) \circ T_s + (\operatorname{div} V(s)) \circ T_s \mathbf{I}] ds$, we get

$$\begin{aligned} -\partial_t \left[\exp \left\{ \int_0^t [{}^*(DV(s)) \circ T_s + (\operatorname{div} V(s)) \circ T_s \mathbf{I}] ds \right\} \theta(t) \right] &= \\ \exp \left\{ \int_0^t [{}^*(DV(s)) \circ T_s + (\operatorname{div} V(s)) \circ T_s \mathbf{I}] ds \right\} F \circ T_t. \end{aligned}$$

By integration we deduce an explicit expression of θ :

$$\begin{aligned}\theta(t) &= \int_t^\tau \exp\left\{-\int_s^t [{}^*DV(\xi) \circ T_\xi + (\operatorname{div}V(\xi)) \circ T_\xi \mathbf{I}] d\xi\right\} F(s) \circ T_s ds \\ &= \int_t^\tau {}^*(DT_t)^{-1} {}^*(DT_s) F(s) \circ T_s \gamma(s) \gamma(t)^{-1} ds\end{aligned}$$

Then taking for $\Lambda = \theta \circ T_t^{-1}$, it is easy to see that Λ is the unique solution of (15)-(16) which, for a suitable right-hand term, will represent the adjoint problem associated to \mathbf{Z} associated to a given cost functional.

□

2.5 A right-hand term supported by $\Sigma(V)$

Let $f \in L^2(\Sigma(V))$ and assume that the mapping $t \mapsto \gamma_{\Gamma_t}^*(f(t)n_t)$ belongs to $L^2((0, \tau), H(D))$. We proved in Theorem 2 the existence of a unique $\Lambda \in C([0, \tau], H(D))$ such that

$$\begin{aligned}-\partial_t \Lambda - D\Lambda.V - {}^*DV.\Lambda - \operatorname{div}V \Lambda &= \gamma_{\Gamma_t}^*(f(t)n_t) \\ \Lambda(\tau) &= 0.\end{aligned}\tag{17}$$

We shall prove that the solution Λ is, in fact, supported by the lateral boundary $\Sigma(V)$ since the right-hand term in this problem is itself supported by $\Sigma(V)$ and there is no diffusion term.

Lemme 5

Let $f \in L^2(0, \tau; L^2(\Gamma_t))$. There exists a unique solution in $C([0, \tau]; L^2(\Gamma_t))$ such that $\partial_t \lambda + \nabla_{\Gamma_t} \lambda.V \in L^2(0, \tau; L^2(\Gamma_t))$ of the following problem

$$\begin{cases} \partial_t \lambda(t) + \nabla_{\Gamma_t} \lambda.V + \lambda \operatorname{div}V &= f(t) \text{ on } \cup_t (\{t\} \times \Gamma_t) \\ \lambda(\tau) &= 0 \text{ on } \Gamma_\tau \end{cases}$$

Proof. Notice that

$$[\partial_t \lambda(t) + \nabla_{\Gamma_t} \lambda.V]_{|\Gamma_t} \circ T_t = \partial_t (\lambda \circ T_t)_{|\Gamma}$$

and consider $\mu \in C([0, \tau]; L^2(\Gamma))$ the unique solution of

$$\begin{aligned}\partial_t \mu + (\operatorname{div}V) \circ T_t \mu &= f(t) \circ T_t \text{ on } (0, \tau) \times \Gamma \\ \mu(\tau) &= 0 \text{ on } \Gamma.\end{aligned}$$

which can be expressed as :

$$\begin{aligned}\mu(t) &= -\int_t^\tau \exp\left\{\int_t^s (\operatorname{div}V) \circ T_r dr\right\} f(s) \circ T_s ds \\ &= -\int_t^\tau \det DT_s.(\det DT_t)^{-1} f(s) \circ T_s ds\end{aligned}$$

Then λ , defined by $\lambda(t) = \mu(t) \circ T_t^{-1}$, is solution of the considered problem. Uniqueness is obvious.

□

Theorem 3

Let $f \in L^2(0, \tau; L^2(\Gamma_t))$. The solution Λ of (17) is supported by $\Sigma(V)$. Precisely

$$\Lambda(t) = -\gamma_{\Gamma_t}^*(\lambda(t)n_t), \quad t \in (0, \tau). \quad (18)$$

where λ is defined in Lemma 5.

For further properties of $p(\cdot)$ see [2] or [5].

Proof. Set $\mathbf{X}(t) = -\gamma_{\Gamma_t}^*(\lambda(t)n_t) (\in H^{-1}(D, R^N))$. We should identify the distribution

$$-\partial_t \mathbf{X} - D\mathbf{X}.V - {}^*DV.\mathbf{X} - (\operatorname{div} V) \mathbf{X}.$$

For that let $\varphi \in \mathcal{D}((0, \tau) \times D)$, thus

$$\begin{aligned} & \langle -\partial_t \mathbf{X} - D\mathbf{X}.V - {}^*DV.\mathbf{X} - (\operatorname{div} V) \mathbf{X}, \varphi \rangle_{\mathcal{D}'((0, \tau) \times D), \mathcal{D}((0, \tau) \times D)} \\ &= - \int_0^\tau \int_{\Gamma_t} \lambda(t) \langle \partial_t \varphi, n_t \rangle d\Gamma_t dt + \int_0^\tau \int_{\Gamma_t} \lambda(t) \langle -D\varphi.V + DV.\varphi, n_t \rangle d\Gamma_t dt \end{aligned}$$

The first term $E_1 = - \int_0^\tau \int_{\Gamma_t} \lambda (\partial_t \varphi).n_t d\Gamma_t dt$ is treated as follows : Using the transformation $T_t(V)$

$$\begin{aligned} E_1 &= - \int_0^\tau \int_{\Gamma} \lambda \circ T_t [(\partial_t \varphi) \circ T_t].n_t \circ T_t \omega(t) d\Gamma dt, \quad \omega(t) = \det(DT_t) \parallel {}^*DT_t^{-1}.n \parallel_{R^N} \\ &= - \int_0^\tau \int_{\Gamma} \langle \partial_t(\varphi \circ T_t) - (D\varphi.V) \circ T_t, n_t \circ T_t \rangle \lambda \circ T_t \omega(t) d\Gamma dt \\ &= \int_0^\tau \int_{\Gamma} \langle \varphi \circ T_t, n_t \circ T_t \rangle \partial_t(\lambda \circ T_t) \omega(t) + \langle \varphi \circ T_t, \partial_t(\omega(t) n_t \circ T_t) \rangle \lambda \circ T_t d\Gamma dt \\ &\quad + \int_0^\tau \int_{\Gamma} \langle (D\varphi.V) \circ T_t, n_t \circ T_t \rangle \lambda \circ T_t \omega(t) d\Gamma dt. \end{aligned}$$

But $\omega(t)n_t \circ T_t = \gamma(t) * DT_t^{-1}n$; where $\gamma(t) = \det DT_t$ so

$$\begin{aligned}
E_1 &= \int_0^\tau \int_\Gamma \partial_t(\lambda \circ T_t) \langle \varphi \circ T_t, n_t \circ T_t \rangle \omega(t) d\Gamma dt \\
&+ \int_0^\tau \int_\Gamma \langle \varphi \circ T_t, \partial_t(\gamma(t) * DT_t^{-1}n) \rangle \lambda \circ T_t d\Gamma dt, \\
&+ \int_0^\tau \int_{\Gamma_t} \lambda \langle D\varphi.V, n_t \rangle d\Gamma_t dt, \\
&= \int_0^\tau \int_\Gamma [\partial_t \lambda + \nabla_{\Gamma_t} \lambda.V] \circ T_t \langle \omega(t) n_t \circ T_t, \varphi \circ T_t \rangle d\Gamma dt \\
&+ \int_0^\tau \int_\Gamma \lambda \circ T_t \langle \varphi \circ T_t, \partial_t(*DT_t^{-1}n) \rangle \gamma(t) d\Gamma dt \\
&+ \int_0^\tau \int_\Gamma \lambda \circ T_t \langle \varphi \circ T_t, \partial_t(\gamma(t)) * DT_t^{-1}n \rangle d\Gamma dt + \int_0^\tau \int_{\Gamma_t} \lambda \langle D\varphi.V, n_t \rangle d\Gamma_t dt \\
&= \int_0^\tau \int_{\Gamma_t} [\partial_t \lambda + \nabla_{\Gamma_t} \lambda.V] \varphi.n_t d\Gamma_t dt + \int_0^\tau \int_{\Gamma_t} \langle D\varphi.V, n_t \rangle \lambda d\Gamma_t dt \\
&+ \int_0^\tau \int_\Gamma \lambda \circ T_t \langle \varphi \circ T_t, n_t \circ T_t \rangle \omega(t) (\operatorname{div} V) \circ T_t d\Gamma dt \\
&- \int_0^\tau \int_\Gamma \lambda \circ T_t \langle \varphi \circ T_t, *(DV) \circ T_t n_t \circ T_t \rangle \omega(t) d\Gamma dt \\
&= \int_0^\tau \int_{\Gamma_t} (\partial_t \lambda + \nabla_{\Gamma_t} \lambda.V) \varphi.n_t d\Gamma_t dt + \int_0^\tau \int_{\Gamma_t} \lambda \langle DV.\varphi, n_t \rangle d\Gamma_t dt \\
&+ \int_0^\tau \int_{\Gamma_t} \lambda \varphi.n_t \operatorname{div} V d\Gamma_t dt - \int_0^\tau \int_{\Gamma_t} \lambda \langle DV.\varphi, n_t \rangle d\Gamma_t dt.
\end{aligned}$$

Finally we obtain

$$\begin{aligned}
&\langle -\partial_t \mathbf{X} - D\mathbf{X}.V - *DV.\mathbf{X} - \operatorname{div} V \mathbf{X}, \varphi_{\mathcal{D}', \mathcal{D}} \rangle \\
&= \int_0^\tau \int_{\Gamma_t} (\partial_t \lambda + \nabla_{\Gamma_t} \lambda.V + \lambda \operatorname{div} V) \varphi.n_t d\Gamma_t dt = \int_0^\tau \int_{\Gamma_t} f(t) \varphi(t).n_t d\Gamma_t dt
\end{aligned}$$

which is equivalent, in a distribution sense, to

$$-\partial_t \mathbf{X} - D\mathbf{X}.V - *DV.\mathbf{X} - \operatorname{div} V \mathbf{X} = \gamma_{\Gamma_t}^*(f(t)n_t)$$

Moreover it is clear that $\mathbf{X}(\tau) = 0$.

From the uniqueness theorem 2, we deduce that $\Lambda(t) = -\gamma_{\Gamma_t}^*(\lambda(t)n_t)$.

□

3 Derivability with respect to the field

As mentioned before, we are interested in the structure of the Eulerian derivative of non-cylindrical functionals of the following type

$$\left\{ \begin{array}{l} j(V) = \int_{Q(V)} F(t, x, u(V)(t, x)) \, dx dt \\ \text{where } Q(V) = \bigcup_{0 < t < \tau} (\{t\} \times T_t(V)(\Omega)) \end{array} \right.$$

τ is a non-negative scalar.

Ω a domain in R^N .

$F : I(= [0, \tau]) \times D \times R^{N'} \rightarrow R \quad C^1$.

The function u is solution of a well-posed non-cylindrical PDE, of order $2m$, $m \in N^*$, in $Q(V)$:

$$\partial_t u + A(u) = f \quad \text{in } Q(V) \tag{19}$$

$$B_j(u) = g_j \quad \text{on } \Sigma(V), \quad 0 \leq j \leq m-1 \tag{20}$$

$$u(0) = u_0 \quad \text{in } \Omega \tag{21}$$

A is a differential operator of order $2m$.

B_j is a boundary differential operator of order m_j ($0 \leq m_j \leq 2m-1$).

In the sequel, the following notations will be used.

Notations 1

Assume the existence of the derivative, at $s = 0$, of the mapping $s \rightarrow u^s(.,.) = u(V + sW)(., T_s(\mathcal{Z}(s)(.)))$ in $L^2(I, H^{2m}(\Omega_t(V)))$ for the weak or the strong topology. It is denoted

$$\dot{u}(V; W).$$

Under the same assumption, let

$$u'(V; W) = \dot{u}(V; W) - \partial_x u \cdot \mathbf{Z}.$$

Lemme 6

Assume that, for any direction $W \in \mathcal{C}([0, \tau], \mathcal{D}(D, R^N))$, the derivative

$\dot{u}(V; W)$ exists in $L^2(I, H^{2m}(\Omega_t(V)))$ and that

$$u'(V; W) \quad \text{depends linearly on } W.$$

Then the functional $j(\cdot)$ is Gâteaux differentiable at V and there exists a time-dependent distribution $G(V) \in L^1(0, \tau; \mathcal{D}'(D, R^N))$ with $\text{spt}[G(V)(t)] \subset \overline{\Omega_t(V)}$ such that

$$j'(V; W) = \int_0^\tau \langle G(V)(t), W(t) \rangle_{\mathcal{D}'(D, R^N), \mathcal{D}(D, R^N)} \, dt.$$

Proof. In the perturbed tube $Q(V + sW)$, the cost functional has the following expression :

$$\begin{aligned} j(V + sW) &= \int_0^\tau \int_{\Omega_t(V + sW)} F(t, x, u(V + sW)(t, x)) dx dt \\ &= \int_0^\tau \int_{T_s(\mathcal{Z}^t)(\Omega_t(V))} F(t, x, u(V + sW)(t, x)) dx dt \\ &= \int_0^\tau \int_{\Omega_t(V)} F(t, T_s(\mathcal{Z}^t)(x), u^s(t, x)) dx dt \end{aligned}$$

Hypothesis *iii*) ensures the existence of

$$\mathbf{j}'(\mathbf{V}; \mathbf{W}) = \frac{d}{ds} \mathbf{j}(\mathbf{V} + s\mathbf{W})|_{s=0}.$$

Precisely we have

$$\begin{aligned} j'(V; W) &= \int_0^\tau \int_{\Omega_t(V)} \partial_x F(t, x, u(t, x)) \cdot \mathbf{Z}(t, x) + \partial_y F(t, x, u(t, x)) \cdot \dot{u}(t, x) \\ &\quad + F(t, x, u(t, x)) \operatorname{div} \mathbf{Z}(t, x) dx dt \end{aligned}$$

which is equivalent to

$$\begin{aligned} j'(V; W) &= \int_0^\tau \langle \partial_y F(t, x, u(t, x)); \dot{u}(t, x) - \partial_x u(t, x) \cdot \mathbf{Z}(t, x) \rangle dt \\ &\quad + \int_0^\tau \int_{\Omega_t(V)} \operatorname{div}[F(t, x, u(t, x)) \mathbf{Z}(t)] dx dt \end{aligned}$$

which we rewrite, under smoothness assumptions, as follows

$$\begin{aligned} j'(V; W) &= \int_0^\tau \int_{\Omega_t(V)} \partial_y F(t, x, u(t, x)) \cdot u'(t, x) dx dt \\ &\quad + \int_0^\tau \int_{\Gamma_t(V)} F(t, x, u(t, x)) \mathbf{Z}(t) \cdot \mathbf{n}_t d\Gamma_t dt. \end{aligned} \tag{22}$$

According to assumption ii) and the linear dependence of \mathbf{Z} on W , we obtain the linear dependence of $\mathbf{j}'(\mathbf{V}; \mathbf{W})$ on W .

□

Considering Λ the solution of problem (15)-(16) we can express the boundary integral on $\mathbf{Z}(t) \cdot \mathbf{n}_t$ explicitly in terms of $W(t) \cdot \mathbf{n}_t$. It is the object of the following lemma.

Lemme 7

Let F be a sufficiently smooth function defined on $\Sigma(V)$. Then

$$\int_0^\tau \int_{\Gamma_t(V)} F(t) \mathbf{Z}(t) \cdot \mathbf{n}(t) d\Gamma_t dt =$$

$$\int_0^\tau \int_{\Gamma_t(V)} \left\{ \int_t^\tau F(s) \circ T_s(V) \circ T_t(V)^{-1} ds \right\} W(t).n(t) d\Gamma_t dt$$

Proof.

$$\begin{aligned} \int_0^\tau \int_{\Gamma_t} F(t) \mathbf{Z}.n_t d\Gamma_t dt &= \int_0^\tau \langle -\partial_t \Lambda - D\Lambda.V - {}^*DV.\Lambda, \mathbf{Z} \rangle dt \\ &= \int_0^\tau \langle \partial_t \mathbf{Z} + D\mathbf{Z}.V - DV.\mathbf{Z}, \Lambda \rangle dt = - \int_0^\tau \int_{\Gamma_t(V)} \lambda(t) W.n_t d\Gamma_t dt \\ &= \int_0^\tau \int_{\Gamma_t} \int_t^\tau F(s) \circ T_s \circ T_t(V)^{-1} W(t).n(t) [\gamma(s)\gamma(t)] \circ T_t(V)^{-1} ds d\Gamma_t dt \end{aligned}$$

□

Remark 3

In fact, we can use the explicite expression of \mathbf{Z} given by (11) to obtain the expression of the integral in terms of W and by-pass the adjoint problem associated to \mathbf{Z} .

In the sequel, it will be shown that the eulerian derivative coincides with the shape derivative when the functional depends only on the shape of the tube. First let us define a *tube function* (resp. *tube functional*).

Lemme 8

If $u(V + W) = u(V)$ (resp. $\mathbf{j}(\mathbf{V} + \mathbf{W}) = \mathbf{j}(\mathbf{V})$), for any sufficiently smooth V and W s.t. $W(t).n_{\Omega_t(V)} = 0$ on $\Sigma(V)$, then u (resp. \mathbf{j}) depends only on the shape of the considered tube. It is called a *tube function* (resp. *tube functional*).

Proposition 5

The hypotheses of Lemma 6 are assumed to be satisfied.

1. If u depends only on the trace, on the lateral boundary $\Sigma(V)$, of the field V , then the gradient $G(V)$ is supported on $\Sigma(V)$ and there exists $R(V) \in L^1(0, \tau; \mathcal{D}'(\Gamma_t(V)))$ s.t.

$$G(V)(t) = \gamma_{\Gamma_t(V)}^*(R(V)(t))$$

where $\gamma_{\Gamma_t(V)}^*$ is the adjoint of the trace operator on Γ_t .

2. If u is a tube function (so it is denoted $u(Q_V)$), then

$$j'(V; W) = 0 \quad \text{for any } W \text{ s.t. } W(t).n_{\Omega_t(V)} = 0 \text{ in } \Gamma_t(V) \text{ for a.e. } t \in [0, \tau].$$

3. Assume Ω of class \mathcal{C}^k , $k \geq 1$ and the linear mapping $W \mapsto \mathbf{j}(\mathbf{V}; \mathbf{W})$ continuous in $\mathcal{C}([0, \tau]; C^k(\overline{D}, R^N))$. Under the assumptions of Lemma 6 and if u is a tube function, there exists a time-dependent distribution g , $g(t) \in [\mathcal{D}^{k-1}(\Gamma_t(V))]'$, such that

$$\mathbf{j}'(\mathbf{V}; \mathbf{W}) = \int_0^\tau \langle \gamma_{\Gamma_t(V)}^*(\mathbf{g}(\mathbf{V})(t)\mathbf{n}_t), \mathbf{W}(t) \rangle_{\mathcal{D}^{k-1}(\mathbf{D})', \mathcal{D}^{k-1}(\mathbf{D})} dt$$

Proof.

1. Let W be a field such that $W(t) \in \mathcal{D}(D, R^N)$ and $\text{spt}W(t) \cap \overline{\Omega_t(V)} = \emptyset$ for any $t \in [0, \tau]$. Thus $T_t(V + sW)\Omega = T_t(V)\Omega (= \Omega_t(V))$. Therefore $Q(V + sW) = Q(V)$. On the other hand, $u(V + sW)(t)|_{\Gamma_t(V)} = u(V)(t)|_{\Gamma_t(V)}$ a.e. in $[0, \tau]$. The well-posedness of the PDE satisfied by u implies its uniqueness. So, in this case, we have $u(V + sW) = u(V)$ a.e. in $Q(V)$. This proves that $\mathbf{j}'(\mathbf{V}; \mathbf{W}) = \mathbf{0}$. So $\text{spt}G(V)(t) \subset \overline{\Omega_t(V)}$ a.e. in $[0, \tau]$. By similar arguments and considering vector fields W such that $W(t) \in \mathcal{D}(\Omega_t(V), R^N)$ for any $t \in [0, \tau]$, we prove that $\text{spt}G(V)(t) \subset \overline{\Omega_t(V)^c}$ for any $t \in [0, \tau]$. Hence, we can conclude that $\text{spt}G(V)(\cdot) \subset \Sigma(V)$.
2. Expressing the functional \mathbf{j} at the “point” $V + sW$, we obtain

$$j(V + sW) = J(V + sW, Q(V + sW)) = J\left(V + sW, \bigcup_{0 < t < \tau} [\{t\} \times T_s^t \Omega_t(V)]\right)$$

where $T_s^t = T_t(V + sW) \circ T_t(V)^{-1}$.

The condition $W(t).n_{\Omega_t(V)} = 0$ for any $t \in [0, \tau]$ implies that $T_t(V + sW)\Omega = T_t(V)\Omega (= \Omega_t(V))$. It comes that $Q(V + sW) = Q(V)$. Moreover $u(V + sW) = u(V)$, then $\mathbf{j}(V + sW) = \mathbf{j}(V)$. Therefore $\mathbf{j}'(\mathbf{V}; \mathbf{W}) = \mathbf{0}$.

3. The continuity of the mapping $W \mapsto \mathbf{j}(\mathbf{V}; \mathbf{W})$ in $\mathcal{C}([0, \tau]; C^k(\overline{D}, R^N))$ and the fact that the gradient $G(V)(t) = \gamma_{\Gamma_t(V)}^*(R(V)(t))$, give that $R(V) \in L^1(0, \tau; \mathcal{D}^{-k+1}(\Gamma_t(V)))$. Moreover

$$\begin{aligned} j'(V; W) &= \int_0^\tau \langle R(V)(t), W(t) \rangle_{\mathcal{D}^{-k+1}, \mathcal{D}^{k-1}} dt \\ &= \int_0^\tau \langle R(V)(t), C(t) \rangle dt + \int_0^\tau \langle R(V)(t), (W(t).n_t(V))n_t(V) \rangle dt \\ &= \int_0^\tau \langle R(V)(t), (W(t).n_t(V))n_t(V) \rangle dt \end{aligned}$$

where C is an admissible vector field such that

$$C(t) = W(t) - (W(t).n_t(V))n_t(V) \text{ on } \Gamma_t(V).$$

We know that $j'(V; C) = 0$. So under the specified hypotheses

$$R(V)(t) = \gamma_{n_t}^*(g(V)(t)n_t(V)).$$

where γ_{n_t} is the normal trace.

□

Remark 4

1. More generally if u depends only on the trace, on the lateral boundary $\Sigma(V)$, of the field V and $R(V) \in L^1(0, \tau; \mathcal{D}^{-k+1}(\Gamma_t(V)))$ then

$$j'(V; W) =$$

$$\int_0^\tau \langle R(V)(t), (W(t).n_t(V)) n_t(V) \rangle dt + \int_0^\tau \langle R(V)(t), C(t) \rangle dt$$

where the first integral of the right-hand term is the shape derivative and the second one is purely dynamic and is due to the variation of the tangential component of W .

2. Assume the existence of a vector function R with $R(t) \in L^p(\Gamma_t(V), R^N)$, $p \geq 1$, such that $G(V)(t) = \gamma_{\Gamma_t(V)}^* R(t)$. Then, under the assumptions of Lemma 6 and if the density g satisfies $g(t) = R(t).n_t \in L^p(\Gamma_t(V))$, we have an integral representation for the derivative

$$j'(V; W) = \int_0^\tau \int_{\Gamma_t(V)} g(V)(t) W(t).n_t d\Gamma_t dt.$$

4 Newton-Shape Method

We apply the previous tube analysis to a usual shape functional $J(\Omega)$. For a given field V , the domain $\Omega_t(V)$ is defined for any t then we get the classical expansion (25) that we treat following the previous tube functional approach. The “initial” domain Ω being given as well as the time t_0 , from (25) the functional $J(\Omega_{t_0}(V))$ turns to be a functional $j(V)$ of the speed vector field V , $j(V) = J(\Omega_{t_0}(V))$ being defined by the right hand side of (25). Then we apply our previous tube derivative calculus with the use of the transverse field Z associated to the perturbation field W . Concerning the shape analysis we adopt now the classical terminologie introduced in [8], [12] concerning the notions of shape derivative (resp. boundary shape derivative) of distributions defined on a moving domain (resp. on a moving boundary or surface). For example the boundary shape derivative of the normal vector field, denoted by $n'_\Gamma(V)$ is given by $n'_\Gamma(V) = -\nabla_\Gamma(< V(0), n >)$ see [5].

Assume the family of domains to be smooth enough, say Ω of class C^k with the integer $k \geq 1$. We know that for any

$$V \in \mathcal{E}_{0,k} \stackrel{def}{=} C^0([0, \tau], \mathcal{V}_o^k(D))$$

the associated flow and its inverse, $T(V)$ and $T(V)^{-1}$, are elements of $\mathcal{C}^1([0, \tau], \mathcal{C}^k(\overline{D}, R^N))$. Consider a shape functional $J(\cdot)$ defined on a such family of domains \mathcal{A} containing the domains

$$\{\Omega_s(Y) = T_s(Y)\Omega; \forall s \in [0, \tau], \forall Y \in \mathcal{V}_o^k(D)\}.$$

Suppose J shape differentiable on any $B \in \mathcal{A}$. We denote by $dJ(B; Y)$ the shape derivative of J on any $B \in \mathcal{A}$ in the direction Y . If B is in $\mathcal{C}^k(\overline{D}, R^N)$, there exists a Distribution

$G(B)$ of finite order, $G(B) \in \mathcal{D}^{k-1}(D, \mathbb{R}^n)'$ such that

$$dJ(B; Y) = \langle G(B), Y \rangle_{\mathcal{D}^{k-1}(D, \mathbb{R}^n)', \mathcal{D}^{k-1}(D, \mathbb{R}^n)}.$$

When $k = 1$ that gradient is a (vector) measure supported by the boundary.

4.1 expansion

We assume that the gradient is smooth enough in time, $(s \rightarrow G(\Omega_s(V))) \in L^1(0, \tau, \mathcal{D}'(D, \mathbb{R}^N))$, then we obtain the following expansion :

$$j(V) = J(\Omega_{t_0}(V)) = J(\Omega) + \int_0^{t_0} \langle G(\Omega_s(V)), V(s) \rangle ds \quad (23)$$

We say that the shape functional J is twice shape differentiable (see [10]) if

- The mapping $\Omega \rightarrow G(\Omega)$ is shape differentiable, that is :
- $\forall V$, the mapping

$$\begin{aligned} R &\longrightarrow \mathcal{D}^{k-1}(D, \mathbb{R}^N)' \\ s &\longmapsto G(\Omega_s(V)) \end{aligned}$$

is derivable. The shape derivative, which is the derivative at $s = 0$ is denoted $G'(\Omega; V)$.

We assume that the mapping $Z \rightarrow G'(\Omega_s(V), Z)$ is continuous and then (see [7], [8]) depends only on $V(s)$ then we get $G'(\Omega_s(V), Z) = G'(\Omega_s(V), Z(s))$.

Moreover we assume that the mapping $Z \rightarrow G'(\Omega_s(V), Z)$ is linear and continuous and since $\partial\Omega_s(V)$ is smooth, from the generic argument developped in [10] we derive that $G'(\Omega, \cdot)$ will depend on the autonomous field V only through its normal component on the boundary, that is $G'(\Omega_s(V)) \cdot Z = \tilde{G}'(\Omega_s(V)) \cdot (z(s))$ where $z(s) = \langle Z(s), n_s \rangle$.

We define the associated *Shape Hessian* (which can be seen as a continuous linear operator) $\tilde{G}'(\Omega_s(V)) \in \mathcal{L}(\mathcal{D}(\Gamma, \mathbb{R}^N), \mathcal{D}(D, \mathbb{R}^N)')$.

- We suppose the mapping $s \mapsto G'(\Omega_s(V))$, to be continuous.
- The derivative of j on V in the direction W is

$$j'(V; W) = \int_0^{t_0} \langle G'(\Omega_s(V)) \cdot \mathbf{Z}, V \rangle + \langle G(\Omega_s(V)), W \rangle ds$$

From [9] the structure theorem ([9] [8], [4]) for gradient, there exists a scalar distribution, called shape density gradient, $g(B)$ in $\mathcal{D}^{k-1}(\partial B)'$ such that

$$dJ(B; Y) = \langle g(B); \gamma_{\partial B}(Y) \cdot n \rangle_{\mathcal{D}^{k-1}(\Gamma)', \mathcal{D}^{k-1}(\Gamma)} \quad (24)$$

We introduce the class of Shape differentiable functionals as follows :

Definition 1

The shape functional J is in $\mathcal{H}^1(\mathcal{A})$ if it is shape differentiable at any $B \in \mathcal{A}$ and if the density gradient $g(B)$ of J at $B \in \mathcal{A}$ verifies : $g(B) \in H^{\frac{1}{2}}(\partial B)$.

Assume J in $\mathcal{H}^2(\mathcal{A})$ and consider the following expansion :

$$J(\Omega_{t_0}(V)) = J(\Omega) + \int_0^{t_0} \langle G(\Omega_s(V)), V(s) \rangle ds \quad (25)$$

We are interested in computing the first order optimality condition for the minimization problem

$$\min_V J(\Omega_{t_0}(V)), \quad t_0 \text{ being fixed.}$$

We suppose that, $\forall V$, the mapping $s \mapsto \tilde{G}'(\Omega_s(V))$, to be continuous. (That is equivalent to say that the mapping $\Omega \rightarrow G'(\Omega)$ is continuous from \mathcal{A} equipped with the *Courant Metric of domains* (see [4]) in $\mathcal{D}(D, R^N)$).

4.2 Second derivative with density gradient formulation

The derivative of j on V in the direction W is

$$j'(V; W) = \int_0^{t_0} \langle G'(\Omega_s(V)) \cdot \mathbf{Z}(s), V \rangle + \langle G(\Omega_s(V)), W \rangle ds$$

If we rewrite j in terms of the density gradient $g(\Gamma_s(V))$ associated to $G(\Omega_s(V))$, using the fact that

$$G(\Omega_s(V)) = \gamma_{\Gamma_s(V)}^*(g(\Gamma_s(V)) n(\Gamma_s(V))) \quad (26)$$

So that

$$j(V) = J(\Omega) + \int_0^{t_0} \int_{\Gamma_s(V)} g(\Gamma_s(V)) \langle V(s), n(\Gamma_s(V)) \rangle d\Gamma_s ds$$

thus the derivative of j at a point V , in the direction W , has a more explicit expression.

Proposition 6

Assume

1. the data sufficiently smooth
2. the shape derivative $g'(\Gamma_s(V); Y)$ exists for any admissible autonomous direction Y and any $s \in [0, t_0]$.

Then $\Gamma_s(V)$,

$$\begin{aligned} j'(V; W) = & \int_0^{t_0} \int_{\Gamma_s(V)} (g'(\Gamma_s(V); \mathbf{Z}(s)) \langle V(s), n(\Gamma_s(V)) \rangle + \\ & g(\Gamma_s(V)) \langle W(s), n(\Gamma_s(V)) \rangle + \operatorname{div}_{\Gamma_s(V)}(g(\Gamma_s(V)) V(s)) \langle \mathbf{Z}(s), n_s \rangle + \\ & \langle \nabla_{\Gamma_s}(g(\Gamma_s(V))) , V_{\Gamma_s(V)}(s) \rangle \langle \mathbf{Z}(s), n_s \rangle \\ & + g(\Gamma_s(V)) \langle DV(s).n_s, n_s \rangle \langle Z(s), n_s \rangle) d\Gamma_s ds \end{aligned}$$

and it depends on $W(s)$ through $Z(s)$.

Proof. By direct calculation, following derivatives rules and notations ([8], [5]), κ being the mean curvature of the manifold, we get :

$$\begin{aligned} j'(V; W) &= \int_0^{t_0} \int_{\Gamma_s(V)} (g'(\Gamma_s(V); \mathbf{Z}(s)) \cdot \langle V(s), n(\Gamma_s(V)) \rangle + \\ &g(\Gamma_s(V)) \cdot \langle W(s), n(\Gamma_s(V)) \rangle + g(\Gamma_s(V)) \cdot \langle V(s), n'(\Gamma_s(v); Z(s)) \rangle + \\ &\kappa g(\Gamma_s(V)) \cdot \langle V(s), n(\Gamma_s(V)) \rangle \cdot \langle \mathbf{Z}(s), n(\Gamma_s(V)) \rangle \\ &+ \langle DV(s) \cdot n_s, n_s \rangle \cdot \langle Z(s), n_s \rangle) d\Gamma_s ds \end{aligned}$$

Now as $n'(\Gamma_s(v); Z(s)) = -\nabla_{\Gamma_s(v)} \langle Z(s), n_s \rangle$, using tangential by part integration we obtain (27). \square

We assume the following linearity : $g'(\Gamma_s(V); \mathbf{Z}(s)) = \tilde{g}'(\Gamma_s(V)) \cdot \langle \mathbf{Z}(s), n_s \rangle$, which is true from the structure derivative theorem ([7], [8]), as soon as $g'(\Gamma_s(V); \mathbf{Z}(s))$ is linear and continuous with respect to Z , together with the regularity of the boundary. The first term can be rewritten as :

$$\begin{aligned} &\int_0^{t_0} \int_{\Gamma_s(V)} (g'(\Gamma_s(V); \mathbf{Z}(s)) \cdot \langle V(s), n(\Gamma_s(V)) \rangle) d\Gamma_s ds \\ &= \int_0^{t_0} \int_{\Gamma_s(V)} (\tilde{g}'(\Gamma_s(V))^* \cdot \langle V(s), n(\Gamma_s(V)) \rangle \cdot \langle \mathbf{Z}(s), n(\Gamma_s(V)) \rangle) d\Gamma_s ds \end{aligned}$$

Let

$$F(s) = \tilde{g}'(\Gamma_s(V))^* \cdot \langle V(s), n(\Gamma_s(V)) \rangle + g(\Gamma_s(V)) \operatorname{div} V(s) + \langle \nabla_{\Gamma_s(V)} g(\Gamma_s(V)), V(s) \rangle$$

Then

$$j'(V; W) = \int_0^{t_0} \int_{\Gamma_s(V)} F(s) \mathbf{Z}(s) \cdot n(s) + g(\Gamma_s(V)) W(s) \cdot n(s) d\Gamma_s ds$$

Finally we derive the result concerning the necessary Optimality condition for $j(V) = J(\Omega_{t_0}(V))$:

Proposition 7

The necessary optimality condition associated to considered minimization problem is :

$$a.e.s \text{ in } (0, t_0) \{ \int_s^{t_0} F(r) \circ T_r(V) \circ T_s(V)^{-1} dr \} + g(\Gamma_s(V)) = 0 \quad a.e. \text{ in } \Gamma_s(V)$$

The proof is based on Lemma 7. We can also use the adjoint problem solution λ and derive explicit expression for the gradient of the functional j in order to derive a second order descent method for the shape functional J . From 23 we get

$$\int_0^\tau \int_{\Gamma_t} F(t) \mathbf{Z} \cdot n_t d\Gamma_t dt = - \int_0^\tau \int_{\Gamma_t(V)} \lambda(t) W \cdot n_t d\Gamma_t dt$$

So that :

$$j'(V; W) = \int_0^{t_0} \int_{\Gamma_s(V)} (\lambda(s) + g(\Gamma_s(V))) \langle W(s), n(s) \rangle d\Gamma ds$$

Where λ solves the backward problem :

$$\partial_t \lambda(t) + \nabla_{\Gamma_t} \lambda \cdot V + \lambda \operatorname{div} V = F(t) \text{ on } \cup_t (\{t\} \times \Gamma_t), \quad \lambda(t_0) = 0 \text{ on } \Gamma_{t_0}$$

4.3 Example

Finally, let us consider the following functional and the associated minimization over V :

$$j(V) = J(\Omega) - \frac{1}{2} \int_0^{t_0} \int_{\Gamma_t(V)} \left(\frac{\partial y_t}{\partial n_t} \right)^2 \langle V(t), n_t \rangle d\Gamma dt$$

The eulerian derivative of j in the direction W is given by :

$$\begin{aligned} j'(V; W) = & -\frac{1}{2} \int_0^{t_0} \int_{\Gamma_t(V)} \left(\left[2 \left(\frac{\partial y_t}{\partial n_t} \right)^2 \langle V(t), n_t \rangle + \operatorname{div} (|\nabla y_t(V)|^2 V) \right] \langle \mathbf{Z}, n_t \rangle \right. \\ & \left. + \left(\frac{\partial y_t}{\partial n_t} \right)^2 \langle W, n_t \rangle \right) d\Gamma dt \end{aligned}$$

So the necessary optimality condition would be the following :

$$\int_t^{t_0} \left[2 \left(\frac{\partial y_s}{\partial n_s} \right)^2 \langle V(s), n_s \rangle + \operatorname{div} (|\nabla y_s(V)|^2 V(s)) \right] \circ T_s(V) \circ T_t^{-1}(V) ds + \left(\frac{\partial y_t}{\partial n_t} \right)^2 = 0$$

Références

- [1] M.-C. Delfour and J.-P. Zolésio. Structure of shape derivatives for non smooth domains. *Journal of Functional Analysis*, 104, 1992.
- [2] M. C. Delfour and J.-P. Zolésio. Shape analysis via oriented distance functions. *J. Funct. Anal.*, 123(1) :1–16, 1994.
- [3] M. C. Delfour and J.-P. Zolésio. *Shapes and Geometries*, volume 4 of *Advances in Design and Control*. SIAM, Philadelphia, August 2001.
- [4] M. C. Delfour and J.-P. Zolésio. *Shapes and geometries*. Advances in Design and Control. Society for Industrial and Applied Mathematics (SIAM), Philadelphia, PA, 2001. Analysis, differential calculus, and optimization.
- [5] F. R. Desaint and J.-P. Zolésio. Manifold derivative in the laplace-beltrami equation. *J. Funct. Anal.*, 151(1) :234–269, 1997.

- [6] R. Dziri and J.-P. Zolésio. Dynamical shape control in non-cylindrical navier-stokes equations. *J. Convexe Anal.*, 1999.
- [7] Jean-Paul and Jean-Paul Zolésio. Identification de domaine par deformation. In *Doctorat d'état en mathématiques*, Université de Nice, pages 1–450. Dekker, New York, 1979.
- [8] Jan Sokółowski and Jean-Paul Zolésio. *Introduction to shape optimization*, volume 16 of *Springer Series in Computational Mathematics*. Springer-Verlag, Berlin, 1992. Shape sensitivity analysis.
- [9] J.-P. Zolésio. *Identification de domaines par déformations*. Thèse de doctorat d'état, Université de Nice, France, 1979.
- [10] J.-P. Zolésio. Introduction to shape optimization problems and free boundary problems. In *Shape optimization and free boundaries (Montreal, PQ, 1990)*, volume 380 of *NATO Adv. Sci. Inst. Ser. C Math. Phys. Sci.*, pages 397–457. Kluwer Acad. Publ., Dordrecht, 1992.
- [11] J.-P. Zolésio. Set weak evolution and transverse field variational applications ans shape differential equation. Rapport de recherche 4649, INRIA, november 2002.
- [12] Jean-Paul Zolésio. The material derivative (or speed) method for shape optimization. In *Optimization of distributed parameter structures, Vol. II (Iowa City, Iowa, 1980)*, volume 50 of *NATO Adv. Study Inst. Ser. E : Appl. Sci.*, pages 1089–1151. Nijhoff, The Hague, 1981.



Unité de recherche INRIA Sophia Antipolis
2004, route des Lucioles - BP 93 - 06902 Sophia Antipolis Cedex (France)
Unité de recherche INRIA Lorraine : LORIA, Technopôle de Nancy-Brabois - Campus scientifique
615, rue du Jardin Botanique - BP 101 - 54602 Villers-lès-Nancy Cedex (France)
Unité de recherche INRIA Rennes : IRISA, Campus universitaire de Beaulieu - 35042 Rennes Cedex (France)
Unité de recherche INRIA Rhône-Alpes : 655, avenue de l'Europe - 38330 Montbonnot-St-Martin (France)
Unité de recherche INRIA Rocquencourt : Domaine de Voluceau - Rocquencourt - BP 105 - 78153 Le Chesnay Cedex (France)

Éditeur
INRIA - Domaine de Voluceau - Rocquencourt, BP 105 - 78153 Le Chesnay Cedex (France)
<http://www.inria.fr>
ISSN 0249-6399